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Transannular Interactions in *Syn-* **and** *Anti-* **[2.2](1,4) Napthalenophanes and** *Syn-* and *Anti-* $\lceil 2.2 \rceil(1,4)$ **Anthracenophanes**

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The UV spectra of the title compounds were interpreted as a result of *PPP-*CI-I calculations of excitation energies and oscillator strengths for the singlet-singlet transitions. A good agreement of the theoretical transitions with the experimental spectrum was found. A strong transannular effect was characteristic for the *ayn* isomers, the effect decreased in the order: *syn-ring* anthracenophane, *syn-ring* naphtalenophane, *anti-ring* naphthalenophane, *anti-ring* anthracenophane. Also the influence of a pseudo-substituent effect was found and discussed.

(Keywords: [2.2] (1,4) Anthracenophanes; CT transitions; Electronic spectra; [2.2] (1,4) Naphtalenophanes; Pseudo-substitucnt effect; Trans-annular interactions)

Transannulare Wechselwirkungen in syn- und anti- $[2.2](1.4)$ Naphthaleno*phanen and syn- und anti- [2.2] (1,4)Anthracenophanen*

Die UV-Spektren der Titelverbindungen wurden mittels *PPP-CI-1-Rech*nungen interpretiert; dabei wurde eine gute Übereinstimmung der theoretischen Übergänge mit den experimentellen Spektren gefunden. Für die syn-Isomeren wurde ein starker transannularer Effekt festgestellt, dessen Einfluß in der Reihe *syn-Ring Anthracenophan, syn-Ring Napthalenophan, anti-Ring* Naphthalenophan, *anti-Ring* Anthracenophan abnimmt. Ebenso wurde ein pseudo-Substituenten-Effekt gefunden und in die Diskussion miteinbezogen.

Introduction

We present a theoretical consideration of the transannular interaction in some [2.2]naphthalenophanes and [2.2]anthraeenophanes, based on theoretical interpretation of their UV spectra. The considered structures were: *syn-ring and anti-ring* $[2.2](1.4)$ naphthalenophanes^{1, 2} and *syn*ring and *anti-ring* $[2.2](1.4)$ anthracenophanes^{3, 4} (Fig. 1). Difficulties⁵⁻⁹ which we already have discussed in papers concerning quinolinophanes¹⁰ and phenyl derivatives of $[2.2]$ paracyclophane¹¹ appeared in connection with calculations of the excitation energy by semiempirical quantum methods. Reason for these difficulties is that the phanes, as layered molecules, do not fulfill all conditions required for the introduction of the *Born-Oppenheimer* approximation. However, if the main objective of the calculations are excitation energies, it is possible to take into account all vibronic effects by adopting a special procedure for the evaluation of the resonance integral.

A critical step is a choice of a proper procedure to evaluate the resonance and *Coulomb* integrals, in particular the integrals between the $2p_z$ orbitals belonging to two different moieties. Following our treatments^{5, 10, 11}, we applied the *Hinze-Beveridge* parametrization¹² in a modification 10 which appeared to be very effective in the calculation of electronic spectra of cyclophanes. The application of a new empirical parameter is fully justified in case of the layered compounds, as documented by *Vogler 13.*

Materials and Methods

The investigated *syn-* and *anti-ring[2.2J(l,4)naphthalenophanes* were obtained according to *Cram* et al.^{1,2}. The UV spectra were recorded for the solutions in cyclohexane on the UV-Vis 5240 Beckmann and Specord UV-Vis C. Zeiss spectrometers.

The values of excitation energy of the cyclophane chromophores were calculated by a *PPP-CI-1* method in *Hinze-Beveridge* parametrization in the same way as previously¹⁰. For the C parameter, see equ. (1) in ¹⁰, we adopted $C = C_{H \rightarrow B}$ for all orbitals belonging to the same aromatic unit, and $C = 5 C_{H \rightarrow B}$ for orbitals of different aromatic units. In the case of 1,4-dimethylnaphtalene, *DMN,* and 1,4-dimethylanthracene, *DMA,* considered as reference molecules of cyclophane moieties, we used $\omega = -10.00 \text{ eV}$ for the valence state ionization potential (VSAP) of the two carbon atoms connected with the $CH₃$ groups, instead of the usually accepted value of $\omega = -11.16 \, \mathrm{eV}$. This was done in order to account for hyperconjugation effects. As a basis for the geometry parameters we used the structure of $syn\text{-ring}$ $[2.2](1.4)$ anthracenophane^{3, 14}. In particular, we adopted the bond distances, angles, interplanar distances, a degree of boattype bending in a[2.2Jparacyelophane-type structural unit, and deviations from planarity of the remaining structural fragments (Fig. 2).

The analysis of the LCAO coefficients leads to a conclusion that each MO of the dimer, C_i , is characterized by a vector representing an in-phase $(+i)$ and out-of-phase $(-i)$ combination of two vectors of the corresponding chromophores

$$
\mathbb{C}_j(+i) = \begin{pmatrix} \mathbb{C}_i^{ch} \\ \mathbb{C}_i^{ch} \end{pmatrix} \qquad \mathbb{C}_k(-i) = \begin{pmatrix} \mathbb{C}_i^{ch} \\ -\mathbb{C}_i^{ch} \end{pmatrix} \qquad (1)
$$

This means that two MO's of a dimer correspond to each NO of a monomer. This results also in a split of a single excitation of a MO of a monomer into four orbital excitations of a dimer. From the point of view of the VB theory each electronic transition may be expressed as a combination of the electronic excitations within the monomerie moiety (LE) and between the two moieties (CT).

Fig. 1. Considered isomers of naphthalenophane and anthraeenophane: *a syn*ring $[2.2](1,4)$ naphthalenophane; *b* anti-ring $[2.2](1,4)$ naphthalenophane; c 8yn-ring [2.2] (1,4)anthraeenophane *; d anti-ring* [2.2] (1,4)anthraeenophane

Results

The recorded spectra of naphthalenophanes and the spectra of anthracenophanes³ on which the calculations were based are presented in Fig. 3 together with the simulated spectra. The results of the calculations for the conformations assumed are presented in Tables 2 and 3. The tables contain absorption maxima and oscillator strengths related to the experimental data and the split of energy levels of the chromophores, see also Table 1. The data refer to the structures b and c shown in Fig. 2.

Discussion

There is a reasonably good agreement between the theoretical and experimental spectra of the *syn-* and *anti-ring* napthalenophanes and for the syn-ring anthracenophane (see Fig. 3). Energies of the electronic

No.	Sym.	calc.	Energy, nm exp.15	calc.	Oscill. strength exp ¹⁵	Type	Characteristics ^a
					Naphthalene		
1	B_{3u}	307.1	297-320	0.050	0.030	α	$52(5,7)$, 44 $(4,6)$, 2 $(3,9)$
$\overline{2}$	B_{2u}	277.6	266-286	0.610	0.562	p	92(5,6), 8(4,7)
3	B_{3u}	207.5	217-221	2.130	11.70	β	54(4,6), 45(5,7)
$\overline{4}$	B_{2u}	196.2	199-21012	1.246		ß,	89(4,7), 8(5,6), 2(3,8)
					Anthracene		
1	B_{2u}	351.8	323-374	0.613	0.799	p	95(7.8), 4(6.9)
$\boldsymbol{2}$	B_{3u}	337.9		0.079		\propto	$52(7,9)$, 41 $(6,8)$, 3 $(5,11)$
3		231.3	246-252	2.401	21.80		
	B_{3u}					β	54(6,8), 42(7,9), 2(3,10)
$\overline{4}$	B_{2u}	198.2	210-221	1.072	1.148	B'	79(6,9), 12(4,8), 4(7,8)

Table 1. *Characteristics of some selected transitions in naphthalene and anthracene*

a Participation of transitions between naphthalene and anthracene MO's.

Fig. 2. Assumed geometry of the *syn*- and *anti-ring* [2.2](1,4)anthracenophanes: *a syn* isomer. $\alpha/2 = 2.5^\circ$; *b syn* isomer, $\alpha/2 = 4.5^\circ$; *c anti* isomer. $\alpha/2 = 0^{\circ}$; *d anti* isomer, $\alpha/2 = 4.5^{\circ}$. For the *syn*- and *anti*-ring [2.2](1,4) naphthalenophanes the same geometry was assumed. The distance between the aromatic planes, $l = 3.09 \text{ Å}^{14}$

transitions were calculated by considering 50 single-excited configurations of the singlet—singlet type. It is obvious, however, that only the energetically lowest transitions are of interest for the interpretation of the experimental UV spectra. Sixteen of the latter, originating from the fourfold splitting of the α , p, β , and β' bands *(Clar* notation) of the naphthalene and anthracene molecules are presented in Tables 2 and 3.

Transannular Interactions

Table 2. Splitting of naphthalene energy levels by the transannular interaction in the syn-ring $\{2.2\}/(1.4)$ naphthalenophane and anti-

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a Participation of the local excited, LE, and charge transfer, CT, transition; b naphthalene, DMN

anthracene, $\bar{D}M\bar{A}$.

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Fig. 3. Comparison of the experimental and theoretical electronic spectra for: $a \, syn\text{-}[2.2](1,4)$ naphthalenophane; $b \, anti\text{-}[2.2](1,4)$ naphtalenophane; $c \, syn[2.2](1,4)$ anthracenophane; $d \, anti\text{-}[2.2](1,4)$ anthracenophane. Theoretical bands are shown in a sequence from the marked band toward the short-wave bands (some bands are out of the presented range). Dotted lines show the experimental spectra¹⁵ of naphthalene (N), 1,4-dimethylnaphthalene (DMN), anthracene (A) and 1,4-dimethylanthracene (DMA)

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Effects of Changes in geometry

Changes in geometry influence the theoretical spectra to a different degree, depending on the *syn-* and *anti-ring* type of isomerism. For the *anti* isomers the calculated electronic transitions are only to a small degree sensitive to a deviation from planarity of both moieties. In a case of an *exo-bending,* the wavelengths change no more than by 2 nm, also the changes of oscillator strengths are very limited. This is not so, however, for the *syn* isomers. A deviation from planarity by 5° and 9° causes the hypsochromic shifts with an average of 8 nm for naphthalenophane and 12 nm for anthracenophane. These changes are the greatest in the longwave region and the smallest in the shortwave region. They amount for naphthalenophane 18 and 2nm, for anthracenophane 22 and 2 nm. The oscillator strengths decrease with increased bending, the changes being more pronounced for naphthalenophane.

Substituent and Pseudosubstituent Effects

Changes in the shortwave range of the experimental spectra (in which the bands for naphthalenophanes are denoted as β and β' and for anthracenophanes as β) correspond qualitatively to the changes which can be found in the spectra of *DMN* and DMA, respectively. This fact points to a specific pseudo-substituent effect originating from the deviation from planarity of the apex carbon atoms localized in the boat-shaped benzene rings.

The substituent effect resulting from the presence of methylene groups is reflected in the splitting and shifting of the theoretical bands. In the spectrum of DMN the band β at 207.5nm and the band β' at 196.2 nm show a splitting into the pairs of bands at 217.0, 199.4nm and 200.2 and 190.0nm. Also, all bands show a bathochromic shift. For DNA the β band at 231.3 nm is split into the 249.0 and 226.8 nm bands. The bands shift bathochromically, except the β' band (Fig. 3).

In the cyclophanes under investigation the pseudo-substituent effect appears only in the shortwave region of the spectra. For naphtalenophanes the effect appears in the bands at 236.5, 180.8, and 177.8nm *(syn-ring)* and 247.3, 191.2nm *(anti-ring);* for anthracenophanes the bands appear at 248.9,224.8 nm and 253.2,240.6 nm. All the examples listed originate from a split of the β bands. The naphthalenophane bands at 236.5 and 247.3nm corresponding to experimental bands at 245.5 *(syn)* and 239.0nm *(anti)* result from a transannular effect. Obviously, this effect does not take place in *DMN,* the spectrum of which shows an absorptions minimum in the considered region. All these examples confirm the existanee of a pseudo-substituent effect. In

the calculations performed neither the π -electrons nor some special parameters were taken into consideration as it is usually done in case of a standard treatment aiming at an interpretation of the usual substituent effect of the methyl groups in *DMN* and *DMA.*

Tran8annular Effects

Napthalenophanes show a weak but still observable absorption in the region near 325nm, which is interpreted as originating from the 301.7 nm band of naphtalene. This proves the presence of a transannular coupling in naphthalenophanes because the band corresponds to the transannular $b_{1\sigma(x)}$ band at 345.9nm, typical for the spectrum of [2.2] paracyclophane⁵. Theoretical considerations and calculations point to the presence of two inactive maxima of absorption in the long-wave range of naphtalenophanes spectra, i.e. above 400 nm. The band for a *syn-ring* isomer shows a bathochromic shift with respect to that of the *anti-ring* isomer. We consider these bands as originated from the 277.6nm band *"p"* of naphtalene. The same effect is observed for anthracenophanes. The longest wavelength bands originate from the ptype band of anthracene at 351.8nm. Both the theoretical and experimental band of the *syn* isomer shows a bathochromic shift not found in the *anti* isomer. This bathochromic shift produces yellow coloration of the *syn* isomer. One can therefore assume that the bathochromic shift in the longwave region describes the extent of the transannular interaction in the phanes under consideration. Such interaction depends on the extent of overlapping of two chromophores and is greater for *syn-ring* isomers than for *anti-ring* isomers. Our calculations confirm this fact both for the case of naphthalenophanes and anthracenophanes. The transannular effect, stronger for the *syn*ring isomers, is reflected in a split of naphthalene and anthracene bands (above 60nm), being greater for the *syn* structures than for the *anti* structures. Accepting the extent of the split as a measure of transannular interactions, one can formulate the following order of decreasing interactions: syn-ring anthracenophane, syn-ring napthalenophane, *anti-ring* naphthalenophane and *anti-ring* anthracenophane.

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